Baseline requirements for detecting biosignatures with the HabEx and LUVOIR mission concepts

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Abstract. A milestone in understanding life in the universe is the detection of biosignature gases in the atmospheres of habitable exoplanets. Future mission concepts under study by the 2020 decadal survey, e.g., Habitable Exoplanet Imaging Mission (HabEx) and the Large UV/Optical/IR Surveyor (LUVOIR), have the potential of achieving this goal. We investigate the baseline requirements for detecting four molecular species, H$_2$O, O$_2$, CH$_4$, and CO$_2$, assuming concentrations of these species equal to that of modern Earth. These molecules are highly relevant to habitability and life on Earth and other planets. Through numerical simulations, we find the minimum requirements of spectral resolution, starlight suppression, and exposure time for detecting biosignature and habitability marker gases. The results are highly dependent on cloud conditions. A low-cloud case is more favorable because of deeper and denser lines whereas a no-cloud case is the pessimistic case for its low albedo. The minimum exposure time for detecting a certain molecule species can vary by a large factor (~10) between the low-cloud case and the no-cloud case. For all cases, we provide baseline requirements for HabEx and LUVOIR. The impact of exozodiacal contamination and thermal background is also discussed and will be included in future studies.

Keywords: Habitable Exoplanet Imaging Mission; Large UV/Optical/IR Surveyor; biosignature; life; high dispersion coronagraphy.
Figure 1 summarizes the state-of-the-art performance of direct imaging instruments or testbeds (solid-lined boxes) and for future space missions (i.e., planets around solar-type stars).

High-dispersion coronagraphy (HDC) is a recent technical development that is designed to bridge the contrast gap. HDC combines high contrast imaging (HCI), a single-mode fiber injection unit (FIU),\(^{10}\) and high-resolution spectroscopy (HRS) to filter out stellar light and extract the planet’s signal. Specifically, HCI suppresses stellar light and spatially separates the planet from its host star. The FIU filters out stellar noise at the planet location since the electric field of a stellar speckle does not couple to the fundamental mode of a single-mode fiber, whereas up to \(\sim 80\%\) of the planet light couples into the fiber. HRS further distinguishes planet signal from stellar signal by its unique spectral features such as absorption lines and radial velocity. Using this three-pronged starlight suppression, HDC can achieve the high sensitivity to study terrestrial planets in the habitable zone.

This paper focuses on the application of HDC on future space telescopes, e.g., HabEx and LUVOIR. The fundamental question we attempt to address here is whether these missions can detect four potential biosignature and habitability marker gases: \(\text{O}_2\), \(\text{CH}_4\), \(\text{H}_2\text{O}\), and \(\text{CO}_2\), and set mission requirements to achieve this goal. The four molecular species that are chosen have a relatively high concentration in the Earth’s atmosphere and are significant signs of life. In addition, these species have been detected in Earthshine and spacecraft observations of our own planet.\(^{11}\) Detecting other biosignature gases, which have much lower concentrations, is expected to be far more challenging for exoplanets.

The paper is organized as follows. In Sec. 2, we briefly describe the four molecular species we investigate in this paper and the model we use to generate the spectra for our HDC simulations. In Sec. 3, we briefly describe our HDC simulation approach. More details can be found in previous work on HDC.\(^ {9,12}\) Results are given in Sec. 4 followed by summary and conclusion in Sec. 5.

## 2 Molecular Species of Interest

### 2.1 Model Description

The spectra of Earth-like exoplanets are generated by an atmospheric chemistry and radiative transfer model.\(^ {3,13,14}\) We calculate the molecular abundance as a function of altitude, controlled by photochemical and disequilibrium chemistry processes.\(^ {3}\) The model atmosphere has been compared with terrestrial measurements in the midlatitude and closely resembles the present-day Earth.\(^ {3}\) We calculate the disk-integrated reflected light spectra with an eighth-order Gaussian integration and \(\delta\)-two-stream approximation. We include the opacities of \(\text{CO}_2\), \(\text{O}_2\), \(\text{H}_2\text{O}\), and \(\text{CH}_4\) and calculate the planetary flux at a spectral resolution of \(R = \lambda / \Delta \lambda = 500,000\), which is high enough to resolve individual spectral lines of the aforementioned species over \(\lambda = 0.5\) to 5 \(\mu\)m.

It is well known that Earth’s disk-averaged spectrum cannot be reproduced by a single surface type or cloud deck.\(^ {15–17}\) Rather, it requires a combination of a cloud-free surface that is poorly reflective (i.e., the ocean), a highly reflective low cloud at \(\sim 4\) km mimicking the cumulus clouds, and a high cloud at \(\sim 12\) km mimicking the cirrus clouds.\(^ {15}\) We follow this treatment and assume the three types of surface or cloud. We assign equal weights to the three components to derive a reasonable account for an “Earth-like” planet. While the purpose is to reproduce Earth’s spectrum, the resulting spectrum is generally consistent with the Earthshine experiments\(^ {15,16}\) as well as EPOXI measurements.\(^ {19}\)

### 2.2 Earth Spectrum by Molecule

We consider four molecular species in an Earth’s atmosphere, \(\text{CO}_2\), \(\text{O}_2\), \(\text{H}_2\text{O}\), and \(\text{CH}_4\). Their spectra at different spectral resolutions are shown in Fig. 2. Observable line density and line depth decrease as spectral resolution decreases, making detecting certain molecules difficult at low spectral resolutions. Conversely, planet signal is dispersed to more pixels at high spectral resolutions. In such a case, molecule detectability is generally limited by detector noise. Since there are a number of zero- or low-noise detectors that potentially allow us to overcome the detector noise limited case,\(^ {18}\) we assume noiseless detectors in the HDC simulations in this paper.

Planet reflection spectrum is highly dependent on cloud condition. Figure 2 shows four cases: high-cloud, low-cloud, no-cloud, and the average of the three. The high-cloud case (orange) has the highest albedo, but line depth and density are the lowest, because the atmospheric column of absorbing gases above the cloud (\(\sim 12\) km) is the smallest. The effect is more notable for \(\text{H}_2\text{O}\) than for \(\text{O}_2\), because the abundance of \(\text{H}_2\text{O}\) drops more than \(\text{O}_2\) above the cloud. Consequently, the high-cloud case has the advantage of high reflection but the disadvantage of low spectral information content. In comparison, for the no-cloud case (red), lines are denser and deeper and the albedo is as high as the high-cloud case, so the low-cloud case represents the best case scenario for HDC observations. Lastly, the no-cloud case has the lowest albedo at \(< 0.1\), because the cloud-free surface is poorly reflective.

Our averaged spectrum binned down to a low spectral resolution (e.g., \(R = 300\)) is generally consistent with Earthshine
For example, the water absorption feature at 0.95 μm has an absorption depth of ∼50% according to Earthshine observation (see Fig. 6 in Ref. 15), consistent with our averaged spectrum (blue in Fig. 2). In comparison, the low-cloud case alone would have a greater continuum albedo and a deeper absorption.

We, however, recognize that the cloud coverage and heights will not be known a priori for any planets to be observed. In order to cover a large parameter space, we consider all three cases for different cloud conditions: the low-cloud case, the high-cloud case, and the no-cloud case. Results for the average-cloud case and any other cases should fall within the results of the three cases, which will be discussed in subsequent sections.

Figure 2 also emphasizes the need of going beyond 1 μm to detect biosignature gases. There are no strong lines of CH₄ and CO₂ below 1 μm. Detecting these two molecules is difficult at optical wavelengths. Even for molecules with strong lines below 1 μm, going beyond this wavelength would allow many more strong lines to be measured, thus increasing the detectability. We limit our HDC simulations below 1.8 μm, beyond which a cryogenic space mission would be required.


3 Simulations of HDC Observations

3.1 Methodology

In this section, we briefly describe the procedure to simulate HDC observations and conduct data reduction. For more details in HDC simulations, refer to Refs. 9 and 12.

The planet and stellar spectra are convolved with a kernel that corresponds to a certain spectral resolution. The stellar signal is reduced by a factor that we refer to as the starlight suppression level. The starlight suppression includes the suppression from the coronagraph, wavefront control, and additional nulling boost provided by the use of a single-mode fiber.

Poisson noise is added to account for photon noise. In addition, noises incurred in data reduction are included, e.g., errors associated with removing stellar light measured by additional fibers in speckle field. One particularly important noise source, speckle chromatic noise, is also taken into account. This noise arises from wavefront control at high starlight suppression levels and prevents us from detecting biosignature gases at low spectral resolutions.

We briefly discuss here how the speckle chromatic noise confuses biosignature detection at low spectral resolutions and our strategy to remove the confusion. We refer readers to Sec. 6.4 in Ref. 9 for more details. The speckle chromatic noise arises due to wavefront control at deep starlight suppression (< ~ 10^-7). Because wavefront control is normally optimized at a certain wavelength, starlight suppression is shallower at wavelengths that deviate from the optimal wavelength. As a result, speckle would have a parabolic spectral feature that sometimes mimics an absorption band of a molecule species. We name this speckle behavior at deep starlight suppression levels as the speckle chromatic noise. The speckle chromatic noise may be difficult to distinguish from absorption features of potential biosignatures, and this noise source therefore represents a potential false-positive biosignature absorption signal.

At moderately high spectral resolutions, the speckle chromatic noise is no longer a concern because simulation has shown that the noise only has low-frequency spectral features. Therefore, our strategy of removing the speckle chromatic noise is straightforward. We apply a high-pass filter to those simulated observed spectra to remove spectral features with frequencies that correspond to R < 100, where R is spectral resolution. In this procedure, both the speckle chromatic noise and the absorption features are removed. Our HDC simulations are no longer sensitive to spectral features with R < 100. However, the simulations are free of the influence of the speckle chromatic noise. This makes direct comparison of our technique incompatible with other works at low resolutions (R < 100) that do not use a high-pass filter to mitigate the speckle chromatic noise.

The processed spectrum is then cross-correlated with a template spectrum for the molecular species of interest and at the same spectral resolution as the observation. We note that the processed spectrum contains absorption features of all molecule species and have noise appropriately incorporated and thus represents a realistic case for HDC observations. The resulting cross-correlation function (CCF) is used for biosignature detection and to access the relative abundance of different molecules that are present in planet’s atmosphere. The key parameters used in HDC simulations are given in Table 1.

3.2 Detection Definition

Each HDC simulation results in a CCF. We repeat the HDC simulation 100 times for each combination of spectral resolution and starlight suppression level. We define significance of detection in the following way. The maximum CCF value within one spectral resolution element to the planet radial velocity is recorded for each run. After 100 runs, the distribution of maximum CCF values is obtained. We then calculate 16 and 84 percentiles, half of the difference between the two values corresponds to 1 - σ of the distribution. This way of defining 1 - σ is robust against outliers. The maximum CCF value divided by 1 - σ gives detection significance. The median of detection significance values for 100 runs is recorded.

4 Results

In this section, we summarize the requirements for detecting certain molecular species in terms of exposure time, spectral resolution, and starlight suppression level. Listed in order of decreasing difficulty, the molecules of interest are: CH₄, CO₂, O₂, and H₂O. The order of O₂ and H₂O may switch depending on spectral resolution and starlight suppression level. Results in this section will provide the baseline requirements for the three cases with different cloud conditions.

4.1 Requirements for HabEx

4.1.1 No-cloud case

The no-cloud case has the lowest albedo (Fig. 2) and thus represents the worst case scenario. CH₄ cannot be detected even

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Table 1 Inputs for HDC simulations for a Sun–Earth system.

<table>
<thead>
<tr>
<th>Telescope/instrument</th>
<th>Star</th>
<th>Planet</th>
</tr>
</thead>
<tbody>
<tr>
<td>Telescope aperture</td>
<td>T&lt;sub&gt;eff&lt;/sub&gt;</td>
<td>5800 K</td>
</tr>
<tr>
<td>End-to-end throughput</td>
<td>log(g)</td>
<td>4.5</td>
</tr>
<tr>
<td>Spectral resolution</td>
<td>V sin i</td>
<td>2.7 km/s</td>
</tr>
<tr>
<td>Exposure time</td>
<td>Orbital inclination</td>
<td>50 deg</td>
</tr>
<tr>
<td>Wavelength</td>
<td>Radial velocity</td>
<td>0.0 km/s</td>
</tr>
<tr>
<td>Detector noise</td>
<td>Distance</td>
<td>5 pc</td>
</tr>
</tbody>
</table>
with an exposure time of 1600 h (as shown in green contours in Fig. 3). Because of the low albedo, photon count is extremely low for the 4-m aperture of HabEx, there is on average 0.4 photon per pixel at the highest spectral resolution that we consider (R = 51,200) for a 100-h integration time. In addition, CH₄ lines are intrinsically rare and shallow at wavelengths shorter than 1.8 μm (see Fig. 2). Detection of CO₂ is not possible even with an exposure time of 1600 h because of the same reason as the CH₄ case for an albedo <0.1.

The minimum spectral resolution for CH₄ detection is at R = 5000 for a 1600-h exposure time and a starlight suppression level of 1 x 10⁻¹⁰. It has been noted that a higher spectral resolution would help to increase detectability. However, even at R = 256,000 and 1 x 10⁻¹⁰ starlight suppression, the minimum exposure time is 850 h.

CO₂ can be detected for two exposure times that we consider, 400 and 1600 h, for which the minimum spectral resolution is R = 8000 and R = 700 at 1 x 10⁻¹⁰ starlight suppression, respectively. At R = 51,200, the starlight suppression requirement can be relaxed to 5 x 10⁻⁹ and 1 x 10⁻⁹ for 1600- and 400-h exposure times.

The requirement for O₂ detection is less stringent than CO₂ detection. Spectral resolution higher than 100 will suffice for starlight suppression levels deeper than 2 x 10⁻¹⁰ and 1 x 10⁻⁹ for exposure time of 400 and 1600 h. The requirement for starlight suppression is further relaxed toward higher spectral resolutions with increasing exposure time (see orange contours in Fig. 3).

4.1.2 High-cloud case

The high-cloud case has a higher albedo but line depth is much shallower than the no-cloud case. However, the high albedo outweighs the shallow line depth, making it possible to detect all four gases.

The minimum spectral resolution for CH₄ detection is at R = 5000 for a 1600-h exposure time and a starlight suppression level of 1 x 10⁻¹⁰. Detection of CO₂ is not possible even with an exposure time of 1600 h because of the same reason as the CH₄ case for an albedo <0.1.

Fig. 3 Exposure time contours (in hours) for 5 – 7σ detection. Adjacent contours differ by a factor of 4 in exposure time. The simulation is for HabEx and assumes three cases with different cloud conditions: the no-cloud case (green), the high-cloud case (orange), and the low-cloud case (red).
Because the speckle chromatic noise is considered in our simulation, any low-frequency spectral features are removed because of a high-pass filter that we have applied in the data reduction procedure. As a result, no molecules can be detected at a spectral resolution lower than $R = 100$.

The minimum exposure time for $\text{H}_2\text{O}$ detection is 12.7 h at $R = 25,600$ and $C = 1 \times 10^{-10}$, where $C$ denotes starlight suppression level. The requirement for starlight suppression can be relaxed by three orders of magnitude (compared to planet-star contrast at $\sim 10^{-10}$) at high spectral resolutions ($R > 25,600$) and for a long exposure time (i.e., 1600 h). In this regime, the sensitivity boost of high dispersion coronagraphy is fully taken advantage because of sufficient photon flux and high spectral resolution.

### 4.1.3 Low-cloud case

The main difference between the low-cloud case and the high-cloud case is line depth and density (see Fig. 2): spectral information is richer for the low-cloud case, which makes the case more amenable for spectroscopic observation. Red contours in Fig. 3 shows $5 - \sigma$ detection significance contour. Notably, the $5 - \sigma$ extends to lower spectral resolutions and shallower starlight suppression levels for $\text{H}_2\text{O}$ and $\text{CH}_4$. Consequently, detection of these two molecule species is easier than the high-cloud case. This can be explained as follows: atmospheric column density for $\text{H}_2\text{O}$ and $\text{CH}_4$ drops more than $\text{CO}_2$ and $\text{O}_2$, so a high cloud affects $\text{H}_2\text{O}$ and $\text{CH}_4$ detectability more than $\text{CO}_2$ and $\text{O}_2$.

For $\text{CH}_4$, minimum exposure time is 300 h at $R = 51,200$ and $C = 1 \times 10^{-10}$. For longer exposure times, the requirement for both spectral resolution and starlight suppression can be relaxed. For example, the minimum spectral resolution is at $R = 1400$ and $R = 5000$ at $C = 1 \times 10^{-10}$ for exposure time of 1600 and 400 h. At a spectral resolution of 51,200, starlight suppression can be relaxed to $3 \times 10^{-10}$ and $2 \times 10^{-9}$ for exposure time of 400 and 1600 h.

Detection of $\text{CH}_4$ in different cloud conditions exemplifies the influence of planetary spectrum on detection. The variety of possible planetary spectra may be the major uncertainty in estimating the minimum requirements for future space missions in search for biosignatures. We show here the influence of cloud condition. There are other uncertainties that may affect the mission sensitivity to biosignatures, for instance, the time evolution of Earth atmosphere. $\text{CH}_4$ concentration was much higher in Archean Earth atmosphere. Therefore, it is much easier to search for $\text{CH}_4$ for an Archean Earth.

The minimum exposure time for $\text{CO}_2$ detection is 150 h at $R = 51,200$ and $C = 1 \times 10^{-10}$. The minimum spectral resolutions are $R = 700$ and $R = 2200$ for exposure time of 1600 and 400 h at a starlight suppression level of $1 \times 10^{-10}$. At the highest

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**Fig. 4** Exposure time contours (in hours) for $5 - \sigma$ detection. Adjacent contours differ by a factor of 4 in exposure time. The simulation is for LUVOIR and assumes three cases with different cloud conditions: the no-cloud case (green), the high-cloud case (orange), and the low-cloud case (red).
spectral resolution, we consider \((R = 51,200)\) the minimum starlight suppression requirements are \(7 \times 10^{-10}\) and \(2 \times 10^{-9}\) for exposure time of 400 and 1600 h.

The requirement for \(O_2\) detection is similar to that for the high-cloud case.

The minimum exposure time for \(H_2O\) detection is 5 h at \(R = 12,800\) and \(C = 1 \times 10^{-10}\). The requirement for starlight suppression can be relaxed by three orders of magnitude (compared to planet-star contrast at \(\sim 10^{-10}\)) at high spectral resolutions \((R > 25,600)\) and for long exposure times (>400 h).

4.2 Requirements for LUVOIR

4.2.1 No-cloud case

Green contours in Fig. 4 show \(5 - \sigma\) detection contours at different exposure times for biosignatures for the LUVOIR no-cloud case. With a bigger aperture and thus higher photon flux, detection significance contours extend to lower spectral resolutions and shallower starlight suppression levels. Notably, the starlight suppression requirements are relaxed to \(\sim 10^{-8}\) and \(2 \times 10^{-9}\) for \(H_2O\) and \(O_2\) for an exposure time of 100 h. We expect 100 h as a reasonable exposure time for LUVOIR because it is a general purposed mission. \(CO_2\) and \(CH_4\) are not detectable given the 100-h exposure time. The minimum exposure time for \(CO_2\) and \(CH_4\) detection is 183 and 354 h at \(R = 25,600\) and \(C = 1 \times 10^{-10}\). At the same combination of spectral resolution and starlight suppression level, the minimum exposure time for \(H_2O\) and \(O_2\) detection is 5.0 and 25.0 h.

4.2.2 Low-cloud case

Red contours in Fig. 4 show \(5 - \sigma\) detection contours at different exposure times for biosignatures for the LUVOIR low-cloud case. This represents the best case scenario for biosignature search among all that we have considered in this paper. With an exposure time that shorter than 25 h, all biosignatures can be detected at \(R = 51,200\) and \(C = 1 \times 10^{-10}\). However, we must point out that the minimum exposure time is very sensitive to what type of cloud a planet may have. The minimum exposure time usually varies by a large factor. For example, the minimum exposure time for \(CH_4\) detection differs by a factor of more than 10 between the no-cloud case (354 h) and the low-cloud case (24 h).

4.2.3 High-cloud case

Orange contours in Fig. 4 shows \(5 - \sigma\) detection contours at different exposure times for biosignatures for the LUVOIR high-cloud case. For, \(H_2O\) and \(CH_4\), the high-cloud contours recede toward higher spectral resolutions and deeper starlight suppression (i.e., upper-left corner of the parameter space in Fig. 4) when compared to the low-cloud case, indicating lower detectability of these two species in the high-cloud case. The high-cloud contours for \(O_2\) are comparable between the low-cloud case and the high-cloud case, which is consistent with the HabEx result. The high-cloud contours for \(CO_2\) indicate a higher detectability than the low-cloud case. However, we will show in Sec. 5.1 that the detectability is in fact comparable within uncertainties in exposure time calculation.

5 Summary and Conclusion

We study the detectability of four molecular species as biosignature or habitability indicators: \(CH_4\), \(CO_2\), \(O_2\), and \(H_2O\). We conduct HDC simulations to set the minimum requirements of exposure time, spectral resolution, and starlight suppression for the detection of these biosignatures. Major findings are summarized in Figs. 3 and 4. The implications are discussed in details in Sec. 4. The results provide a baseline for mission design in order to search for biosignature gases and study the habitability of exoplanets.

The results of this paper are based on HDC simulations that focus on performance at high spectral resolutions. While the simulations are compatible with low-resolution results, the cross-correlation may not be the best possible way to detect planets or molecules. A more straightforward way is to conduct a conventional ADI/SDI sequence, detecting the planet, obtaining a low-resolution spectrum, and inferring the molecular presence by measuring the absorption band depth. However, we provide two arguments for using the cross-correlation technique at low spectral resolutions. First, this ensures consistency in comparing performance over a broad range of spectral resolutions. Second, speckle chromatic noise may hinder molecule detection at low spectral resolutions with conventional method, the treatment in our HDC simulation ensures the effect of the speckle chromatic noise is properly modeled and removed.

5.1 Comparison to Previous Results

One major difference between this work and previous works on observing Earth-like planet using low-resolution spectroscopy:90 we consider speckle chromatic noise, which is a significant noise source at low spectral resolution (e.g., \(R < 100\)).

Another difference is the definition of detection significance. In this work, we adopt the local variation of CCF peaks rather than the variation estimated from other parts of CCF. In previous works, the variation was estimated using the fluctuation of first quarter and/or the fourth quarter of CCF. In principle, there should be no significant difference. However, using the local variation is more representative of CCF variation around CCF peak.

We compare our results to previous HDC simulations for HabEx and LUVOIR.9,21 Using the same input planet spectrum, i.e., the average of the no-cloud case and the low-cloud case, we compare the results for HabEx 400-h exposure and LUVOIR 100-h exposure (Fig. 5) to Figs. 20 and 15 in Ref. 9 to quantify the influence of difference definition of detection significance. We find that the contours in these two works track each other well with a maximum horizontal (i.e., starlight suppression level) difference by a factor of 2 at spectral resolution higher than \(R \sim 2000\). Part of the difference arises from statistical noise as it is evident that the contours are not smooth. The major difference between results in this paper and Figs. 20 and 15 in Ref. 9 is that we include speckle chromatic noise, which affects results at low spectral resolution significantly. Therefore, it is expected that contours in Fig. 5 deviate from contours in Figs. 20 and 15 in Ref. 9 at lower spectral resolutions.

We also compare our results to Table 2 in Ref. 21 to check the consistency of exposure time calculation for different definitions of detection significance. We can check only one HabEx case for starlight suppression level shallower than \(10^{-10}\) because HDC simulations in this work do not go deeper than \(10^{-10}\). At \(R = 6400\) and \(C = 5 \times 10^{-10}\), the exposure time for \(O_2\) and \(H_2O\) detection is 162 and 404 h. These values are comparable with values in Table 2 in Ref. 21 within a factor or 1.6. We also compare results for two LUVOIR cases (\(R = 6400\) and \(R = 25,600\))
at $C = 1 \times 10^{-9}$). For $R = 6400$ and $C = 1 \times 10^{-9}$, the exposure time for CH$_4$, CO$_2$, O$_2$, and H$_2$O detection is >1600, 406.2, 72.1, and 25.5 h. Except for CH$_4$, exposure times for all other species are within a factor of 1.4 when compared to values in Table 2 in Ref. 21. For CH$_4$, we have a lower limit for exposure time because our simulation stops at 1600 h. For $R = 25,600$ and $C = 1 \times 10^{-9}$, the exposure time for CO$_2$, O$_2$, and H$_2$O detection is 196.0, 39.0, and 18.0 h. These values are also within a factor of 2 to values in Ref. 21.

Given the comparison with previous works, we conclude that the uncertainty of exposure time due to definition of detection significance and Poisson statistical noise is at most a factor of 2. This uncertainty is smaller than the uncertainty due to cloud condition, which can change exposure time by a factor of ∼10 between the low-cloud case and the no-cloud case. To account for the uncertainty of different definition of detection significance, we advise readers to use an error bar of 0.3 dex (i.e., a factor of 2) when using the minimum exposure time estimated in this paper for a given case of cloud condition.

For the speckle chromatic noise, we use a high-pass filter to remove speckle chromatic noise. As a result, low-frequency spectral features are also removed. This results in a much longer exposure time at low spectral resolution to reach a certain detection significance. For example, the exposure time is 80 h for $R = 100$ and $C = 1 \times 10^{-10}$ for O$_2$ detection. In comparison, Ref. 20 estimated the exposure time is 200 h for $R = 70$ and $C = 1 \times 10^{-10}$. The difference can be explained as follows. First, the distance of the planet in their calculation was 10 pc. Factoring the difference would reduce the exposure time from 200 to 50 h. Second, our error bar is 0.3 dex, so the actual exposure time can vary from 40 to 160 h depending on the definition of detection significance. Therefore, the two results are consistent within uncertainty.

5.2 Exozodiacal Flux and Thermal Background

We consider in our HDC simulation only starlight suppression levels that expected to be achieved by a coronagraphic system. However, the flux from exozodiacal dust may set a floor for the achievable starlight suppression level. This floor is estimated at a few times the planet-star contrast for a 4-m aperture, i.e., $\sim 10^{-10}$. LUVOIR designs are less prone to exozodiacal flux because of a smaller point spread function and thus smaller contamination area. This is a result of the larger aperture of LUVOIR.

Another noise source that is not accounted for in our HDC simulation is the thermal background emission from the telescope and instrument. This is not a typical concern at optical wavelengths. However, thermal background quickly becomes a severe issue at longer wavelengths. It is possible to consider a cryogenic system that cools the instrument to minimize thermal background and allows the instrument to operate at longer wavelengths.
wavelengths. For example, mission designs exist that cool the system to 130 K without active cryogenic cooling: Euclid uses a passive cooling radiator system to cool the system to 130 K.\textsuperscript{24} To be more realistic, we will incorporate both exozodiacal light emission and thermal noise in future versions of HDC simulation packages.

5.3 Inner Working Angle in Near Infrared Observation

While observing in near-infrared wavelengths probes biosignatures, such as CH\textsubscript{4}, which is otherwise undetected in optical wavelengths, there are challenges for coronagraphic observations: the inner working angle (in unit of \(\lambda/D\)) decreases due to increasing wavelengths for a given planet-star separation. For example, an Earth–Sun system at 10 pc has an angular separation of 0.1\textsuperscript{\circ}. In comparison, 1 \(\lambda/D\) for a 4-m HabEx at 1.7 \(\mu\)m, where CH\textsubscript{4} features are, is 0.087\textsuperscript{\circ}. While extremely challenging, new concepts have been proposed to address coronagraphy at sub-\(\lambda/D\), e.g., the Single-mode Complex Amplitude Refinement (SCAR) coronagraph,\textsuperscript{25,26} and the vortex fiber nuller (VFN). Moreover, given the long exposure time that is required to search for biosignatures for a 5-pc system, observing nearby systems much closer than 10 pc will not only keep the exposure time tractable but also alleviate the inner working angle problem.

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References


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